# Investigation of the Local Lattice Structure and the Effects of the Orbital Reduction Factor on the g Factors of a Trigonal $[Ni(H_2O)_6]^{2+}$ Cluster in $NiTiF_6 \cdot 6H_2O$ and $ZnSiF_6 \cdot 6H_2O$ Crystals at Different Temperatures

Hui-Fang Li<sup>a</sup>, Xiao-Yu Kuang<sup>a,b</sup>, Huai-Qian Wang<sup>a</sup>, and Ai-Jie Mao<sup>a</sup>

Reprint requests to H.-F. L.; E-mail: scu\_kxy@163.com

Z. Naturforsch. 63a, 609 - 615 (2008); received March 10, 2008

The local octahedral environment of  $\mathrm{Ni}^{2+}$  in  $\mathrm{NiTiF_6} \cdot 6\mathrm{H_2O}$  and  $\mathrm{ZnSiF_6} \cdot 6\mathrm{H_2O}$  crystals with a trigonal distortion has been studied at different temperatures, based on the complete energy matrices. The calculated results showed that the local lattice structure around an octahedral  $\mathrm{Ni}^{2+}$  centre in  $\mathrm{NiTiF_6} \cdot 6\mathrm{H_2O}$  and  $\mathrm{ZnSiF_6} \cdot 6\mathrm{H_2O}$  exhibits a compression distortion. Simultaneously, the orbital reduction effect on the g factors has been studied. The relationship between  $\Delta g = g_{\parallel} - g_{\perp}$  and orbital reduction factor k at 4.2, 77 and 298 (302) K has been discussed, suggesting that there is an almost linear relation between k and k0 for the k1 ion in k2 ion in k3 NiTiF6 · 6H2O and k4 and k5 at each temperature.

Key words: Local Structure; Orbital Reduction Effect; EPR Spectrum; Complete Energy Matrices.

#### 1. Introduction

Nickel fluorotitanate hexahydrate (NiTiF<sub>6</sub>·6H<sub>2</sub>O) and zinc fluorosilicate hexahydrate (ZnSiF $_6 \cdot 6H_2O$ ) belong to the series of compounds having the general formula  $M(II)X(IV)F_6 \cdot 6H_2O$  with M = Zn, Ni, Mg, Co, Fe and X = Si, Ti, Sn, Zr. The fluorosilicates and fluorotitanates have a trigonally distorted CsCltype structure including the two complex ions: the hydrated metal complex  $[M(H_2O)]^{2+}$  and the silicon (or titanium) hexafluoride ion [XF<sub>6</sub>]<sup>2-</sup>. Since Pauling [1] has determined the structure in the space group  $R\bar{3}$ for NiSnCl<sub>6</sub> · 6H<sub>2</sub>O, this local structure symmetry was assumed to be typical of the fluorosilicate and fluorotitanate hexahydrate series. Many physicists and chemists have been devoted to studying these compounds. For instance, the zero-field splitting (ZFS) parameter D and g factors  $(g_{\parallel}, g_{\perp})$  for NiTiF<sub>6</sub>·6H<sub>2</sub>O crystals at different temperatures have been reported by Rubins et al. [2]. The first EPR studies of pair spectra were made on Ni<sup>2+</sup> pairs in ZnSiF<sub>6</sub> · 6H<sub>2</sub>O by Al'tshuler and Valishev [3]. These results provide important information to further studies of the transition metal Ni<sup>2+</sup> ion. The EPR parameters and the temperature dependence of D for Ni<sup>2+</sup> in ZnSiF<sub>6</sub> · 6H<sub>2</sub>O crystals were measured [4-7]. The magnitude of  $\partial D/\partial T$ 

has a maximum of  $(0.0027 \pm 0.0001)$  cm<sup>-1</sup>/K, decreasing to about 0.0015 cm<sup>-1</sup>/K at room temperature [5]. However, no theoretical considerations of the effect of temperature on the local structure and the orbital reduction effect on the g factors for Ni<sup>2+</sup> in NiTiF<sub>6</sub> · 6H<sub>2</sub>O and ZnSiF<sub>6</sub> · 6H<sub>2</sub>O crystals have been performed. In this paper, the local structure parameters R and  $\theta$  at different temperatures are determined by simulating the EPR and optical spectra based on the complete energy matrices for a d<sup>8</sup> ion in a trigonal ligand-field. Meanwhile, by considering the orbital reduction effect, the relationship between the orbital reduction factor k and  $\Delta g$  has been discussed.

### 2. Theory

The spin Hamiltonian, including the ZFS and Zeeman terms, can be written as [8]

$$\hat{H}_{S} = g_{\parallel} \beta H_{z} S_{z} + \beta g_{\perp} (H_{x} S_{x} + H_{y} S_{y}) + D \left( S_{z}^{2} - \frac{1}{3} S(S+1) \right).$$
 (1)

From the spin Hamiltonian, the ZFS parameter *D* can be calculated by the splitting energy levels in the

0932-0784 / 08 / 0900-0609 \$ 06.00 © 2008 Verlag der Zeitschrift für Naturforschung, Tübingen · http://znaturforsch.com

<sup>&</sup>lt;sup>a</sup> Institute of Atomic and Molecular Physics, Sichuan University, Chengdu 610065, China

<sup>&</sup>lt;sup>b</sup> International Centre for Materials Physics, Academia Sinica, Shenyang 110016, China

ground state <sup>3</sup>A<sub>2</sub> for a zero magnetic field:

$$E(\pm 1) = \frac{1}{3}D, \quad E(0) = -\frac{2}{3}D,$$
  
 $\Delta E = E(\pm 1) - E(0) = D.$  (2)

For Ni<sup>2+</sup> in NiTiF<sub>6</sub> · 6H<sub>2</sub>O and ZnSiF<sub>6</sub> · 6H<sub>2</sub>O the structure symmetry of the water octahedron surrounding Ni<sup>2+</sup> belongs to the  $R\bar{3}$  space group. The perturbation Hamiltonian for a d<sup>8</sup> Ni<sup>2+</sup> ion in this system can be written as [9]

$$\hat{H} = \hat{H}_{ee} + \hat{H}_{SO} + \hat{H}_{LF} + \hat{H}_{Zeeman}$$

$$= \sum_{i < j} \frac{e^2}{r_{i,j}} + \zeta \sum_i l_i s_i + \sum_i V_i$$

$$+ \sum_i \mu_B(k \vec{l}_i + g_e \vec{s}_i) \cdot \vec{H},$$
(3)

where  $\hat{H}_{ee}$  is the electron-electron repulsion interaction,  $\hat{H}_{SO}$  is the spin-orbit coupling interaction,  $\hat{H}_{LF}$  is the ligand-field interaction,  $\hat{H}_{Zeeman}$  is the Zeeman interaction,  $V_i$  is the ligand-field potential, and k (0.5 <  $k \le 1$ ) is the orbital reduction factor [10]. The ligand-field potential  $V_i$  may be expressed as

$$V_{i} = \gamma_{00}Z_{00} + \gamma_{20}r_{i}^{2}Z_{20}(\theta_{i}, \varphi_{i}) + \gamma_{40}r_{i}^{4}Z_{40}(\theta_{i}, \varphi_{i}) + \gamma_{23}^{c}r_{i}^{4}Z_{43}(\theta_{i}, \varphi_{i}) + \gamma_{23}^{c}r_{i}^{4}Z_{43}^{5}(\theta_{i}, \varphi_{i}),$$
(4)

where  $r_i$ ,  $\theta_i$  and  $\varphi_i$  are spherical coordinates of the i-th electron. The Zeeman operator can be written by the parallel or perpendicular component to the  $C_3$  axis as follows:

$$\hat{H}_{\parallel} = \sum_{i} \mu_{\mathrm{B}}(k\hat{l}_{iz} + g_{\mathrm{e}}\hat{s}_{iz}) \cdot H_{z},\tag{5}$$

$$\hat{H}_{\perp} = \sum_{i} \mu_{\rm B}(k \hat{l}_{ix} + g_{\rm e} \hat{s}_{ix}) \cdot H_{x}. \tag{6}$$

Then, from (3) we establish the complete energy matrices in the trigonal ligand-field. The matrix elements can been expressed as functions of the Racah parameters B and C, the spin-orbit coupling coefficient  $\zeta$  and the ligand-field parameters  $B_{20}$ ,  $B_{40}$ ,  $B_{43}^c$ ,  $B_{43}^s$ . Generally, the z-axis is chosen along the threefold axis. However, based on the point charge and superposition model, the ligand-field parameter  $B_{43}^s$  will vanish and  $B_{20}$ ,  $B_{40}$ ,  $B_{43}^c$  can be derived as [11]

$$\begin{split} B_{20} &= \frac{1}{2} \sum_{\tau} G_2(\tau) (3\cos^2 \theta_{\tau} - 1), \\ B_{40} &= \frac{1}{8} \sum_{\tau} G_4(\tau) (35\cos^4 \theta_{\tau} - 30\cos^2 \theta_{\tau} + 3), \ (7) \\ B_{43}^c &= \frac{\sqrt{35}}{4} \sum_{\tau} G_4(\tau) (\cos \theta_{\tau} \sin^3 \theta_{\tau} \cos 3\phi_{\tau}), \end{split}$$

where  $G_2(\tau)$  and  $G_4(\tau)$  are written as

$$G_{2}(\tau) = -q_{\tau}eG^{2}(\tau),$$

$$G_{4}(\tau) = -q_{\tau}eG^{4}(\tau),$$

$$G^{k}(\tau) = \int_{0}^{R_{\tau}} R_{3d}^{2}(r)r^{2}\frac{r^{k}}{R_{\tau}^{k+1}}dr$$

$$+ \int_{R_{\tau}}^{\infty} R_{3d}^{2}(r)r^{2}\frac{R_{\tau}^{k}}{r^{k+1}}dr.$$
(8)

 $\tau$  and  $q_{\tau}$  represent the  $\tau$ -th ligand and its effective charge, respectively,  $R_{\tau}$  in (8) represents the Ni-H<sub>2</sub>O bond length and  $\theta_{\tau}$  in (7) the angle between the Ni-H<sub>2</sub>O bond and the  $C_3$  axis. According to the Van Vleck approximation for the  $G^k(\tau)$  integral [12], we obtain the relations

$$G_2(\tau) = \frac{A_2}{R_{\tau}^3}, \quad G_4(\tau) = \frac{A_4}{R_{\tau}^5},$$
 (9)

where  $A_4 = -eq_\tau \langle r^4 \rangle$ ,  $A_2 = -eq_\tau \langle r^2 \rangle$ ,  $A_2/A_4 = \langle r^2 \rangle / \langle r^4 \rangle$ . In the previous works, the radial wave function of Ni<sup>2+</sup> has been given [13], and the crystal structure and optical spectra of NiSiF<sub>6</sub> · 6H<sub>2</sub>O have been reported [14, 15], from which we can derive  $\langle r^2 \rangle / \langle r^4 \rangle = 0.141029$ ,  $A_4 = 20.9$  a. u. and  $A_2 = 2.9475$  a. u. for the [Ni(H<sub>2</sub>O)<sub>6</sub>]<sup>2+</sup> cluster. According to Curie et al.'s covalence theory [16], the covalence factor N can be used to describe the Racah parameters B and C and the spinorbit coupling coefficient C as follows:

$$B = N^4 B_0$$
,  $C = N^4 C_0$ ,  $\zeta = N^2 \zeta_0$ , (10)

where  $B_0 = 1084 \text{ cm}^{-1}$ ,  $C_0 = 4831 \text{ cm}^{-1}$ ,  $\zeta_0 = 649 \text{ cm}^{-1}$  are the free-ion parameters [17]. Then, by diagonalizing the complete energy matrices, the interrelation between electronic and molecular structure may be established, and the local structure of the octahedral  $[\text{Ni}(\text{H}_2\text{O})_6]^{2+}$  cluster in  $\text{NiTiF}_6 \cdot 6\text{H}_2\text{O}$  and  $\text{ZnSiF}_6 \cdot 6\text{H}_2\text{O}$  crystals can be determined by analyzing the EPR and optical spectra.

## 3. Calculations of the Local Structure of Ni<sup>2+</sup> in NiTiF<sub>6</sub> · 6H<sub>2</sub>O and ZnSiF<sub>6</sub> · 6H<sub>2</sub>O

The NiTiF<sub>6</sub> · 6H<sub>2</sub>O and ZnSiF<sub>6</sub> · 6H<sub>2</sub>O crystals are two of the series of isomorphous compounds. The Ni<sup>2+</sup> (Zn<sup>2+</sup>) ion is surrounded by six water molecules. When considering the Ni<sup>2+</sup> ion in NiTiF<sub>6</sub> · 6H<sub>2</sub>O and ZnSiF<sub>6</sub> · 6H<sub>2</sub>O crystals, the local lattice structure of the octahedral Ni<sup>2+</sup> centre is similar to that of

T(K)	R (Å)	$\Delta  heta$ (°)	$-D  (\mathrm{cm}^{-1})$	$g_{\parallel}$	$g_{\perp}$
4.2	2.0325	0.3945	1.2407	2.2590	2.2489
		0.4045	1.2715	2.2592	2.2488
		0.4144	1.3020	2.2593	2.2487
		0.4245	1.3331	2.2595	2.2486
		0.4345	1.3639	2.2597	2.2486
Expt. [2]			$\textbf{1.3} \pm \textbf{0.01}$	$\textbf{2.27} \pm \textbf{0.02}$	$\textbf{2.25} \pm \textbf{0.02}$
77	2.0338	0.4045	1.2775	2.2599	2.2495
		0.4145	1.3085	2.2601	2.2495
		0.4238	1.3372	2.2603	2.2494
		0.4345	1.3703	2.2604	2.2493
		0.4445	1.4012	2.2606	2.2492
Expt. [2]			$\boldsymbol{1.337 \pm 0.006}$	$\textbf{2.27} \pm \textbf{0.02}$	$\textbf{2.26} \pm \textbf{0.02}$
298	2.0543	0.3745	1.2741	2.2718	2.2616
		0.3845	1.3075	2.2720	2.2615
		0.3854	1.3105	2.2720	2.2614
		0.3945	1.3409	2.2722	2.2614
		0.4045	1.3743	2.2724	2.2613
Expt. [2]			$\textbf{1.31} \pm \textbf{0.01}$	$\textbf{2.28} \pm \textbf{0.02}$	$\textbf{2.27} \pm \textbf{0.02}$

Table 1. The ZFS parameters D and  $g_{\parallel}$ ,  $g_{\perp}$  for NiTiF<sub>6</sub> · 6H<sub>2</sub>O:Ni<sup>2+</sup> at 4.2, 77 and 298 K.

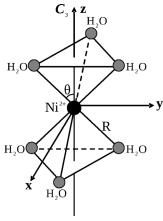


Fig. 1. The local structure of a Ni<sup>2+</sup> centre in NiTiF<sub>6</sub> · 6H<sub>2</sub>O and ZnSiF<sub>6</sub> · 6H<sub>2</sub>O. The Ni<sup>2+</sup> ion is surrounded by six water molecules. R is the Ni-H<sub>2</sub>O bond length and  $\theta$  is the angle between the Ni-H<sub>2</sub>O bond and the  $C_3$  axis.

NiSiF<sub>6</sub> · 6H<sub>2</sub>O. As for the ZnSiF<sub>6</sub> · 6H<sub>2</sub>O:Ni<sup>2+</sup> system, the Ni<sup>2+</sup> ion will occupy the Zn<sup>2+</sup> site, that is surrounded by six neighbour H<sub>2</sub>O molecules. The local structure symmetry around Ni<sup>2+</sup> in NiTiF<sub>6</sub> · 6H<sub>2</sub>O and ZnSiF<sub>6</sub> · 6H<sub>2</sub>O crystals belongs to the  $R\bar{3}$  space group. The Ni<sup>2+</sup> ion is surrounded by six water molecules which make an octahedron along the  $C_3$  axis. In order to describe the local lattice structure of the octahedral [Ni(H<sub>2</sub>O)<sub>6</sub>]<sup>2+</sup> cluster in NiTiF<sub>6</sub> · 6H<sub>2</sub>O and ZnSiF<sub>6</sub> · 6H<sub>2</sub>O crystals, we introduce the two local structure parameters R and  $\theta$ , which denote the Ni-H<sub>2</sub>O bond length and the angle between Ni-H<sub>2</sub>O bond and  $C_3$  axis, respectively, as plotted in Figure 1. As for the NiTiF<sub>6</sub> · 6H<sub>2</sub>O:Ni<sup>2+</sup>

and  $\operatorname{ZnSiF_6} \cdot \operatorname{6H_2O:Ni^{2+}}$  systems, they have a similar crystal structure and the same type of ligand in  $\operatorname{NiSiF_6} \cdot \operatorname{6H_2O}$ , so it is reasonable to approximately apply the optical spectra of  $\operatorname{NiSiF_6} \cdot \operatorname{6H_2O}$  to our calculation. In this case we obtain R = 2.0325 Å, 2.0338 Å, and 2.0543 Å for  $\operatorname{NiTiF_6} \cdot \operatorname{6H_2O}$  and R = 2.0327 Å, 2.0341 Å, and 2.0481 Å for  $\operatorname{ZnSiF_6} \cdot \operatorname{6H_2O:Ni^{2+}}$  at 4.2, 77 and 298 (302) K. Meanwhile, N = 0.969 is obtained and  $k \approx N^2$  can also be evaluated. The local structure parameter  $\theta$  for the octahedral  $[\operatorname{Ni(H_2O)_6}]^{2+}$  cluster in  $\operatorname{NiTiF_6} \cdot \operatorname{6H_2O}$  and  $\operatorname{ZnSiF_6} \cdot \operatorname{6H_2O}$  can be expressed as

$$\theta = \theta_{\rm oh} + \Delta\theta,\tag{11}$$

where  $\theta_{\rm oh}$  denotes the bond angle between the Ni-H<sub>2</sub>O (Zn-H<sub>2</sub>O) bond and the  $C_3$  axis of NiTiF<sub>6</sub>·6H<sub>2</sub>O  $(ZnSiF_6 \cdot 6H_2O)$  in the cubic symmetry  $[\theta_{oh} =$  $\cos^{-1}(1/\sqrt{3}) \approx 54.7356^{\circ}$  [18].  $\Delta\theta$  represents trigonal distortion. Then, in the complete energy matrices, the trigonal distortion angle  $\Delta\theta$  is the only adjustable parameter for Ni<sup>2+</sup> in NiTiF<sub>6</sub> · 6H<sub>2</sub>O and ZnSiF<sub>6</sub> · 6H<sub>2</sub>O:Ni<sup>2+</sup> systems. It can be obtained by fitting the calculated zero-field splitting parameter D and g factors to the observed values. The results are listed in Tables 1 and 2. From our calculation, the EPR parameters (in particular the ZFS parameter D) depend strongly on the change of the bond angle and are not sensitive to the small change of the local structure parameter R. Simultaneously, the optical spectra of  $Ni^{2+}$ in NiTiF<sub>6</sub> · 6H<sub>2</sub>O and ZnSiF<sub>6</sub> · 6H<sub>2</sub>O are listed in Table 3; they are close to each other.

From Tables 1 – 3, by simulating the EPR and optical spectra, the values of R = 2.0325 Å,  $\Delta \theta = 0.4144^{\circ}$ ;

T (K)	R (Å)	$\Delta  heta$ (°)	$-D  (\mathrm{cm}^{-1})$	$g_{\parallel}$	$g_{\perp}$	
4.2	2.0327	0.0244	0.0781	2.2328	2.2325	
		0.0344	0.1100	2.2329	2.2317	
		0.0414	0.1324	2.2330	2.2321	
		0.0544	0.1739	2.2332	2.2321	
Expt.			$0.1325 \pm 0.001$ [5]	$2.233 \pm 0.002$ [5]		
•			$0.129 \pm 0.001$ [6]	2.25±0.01 [6]	$2.24 \pm 0.01*$	
77	2.0341	0.0444	0.1427	2.2338	2.2329	
		0.0544	0.1748	2.2340	2.2327	
		0.0593	0.1905	2.2341	2.2326	
		0.0644	0.2068	2.2342	2.2327	
Expt.			<b>0.1905</b> ± <b>0.0015</b> [5]	$2.235 \pm 0.005$ [5]		
			<b>0.2</b> ± <b>0.01</b> [6]	<b>2.25</b> ± <b>0.01</b> [6]		
			<b>0.191</b> ± <b>0.002</b> [7]	2.23±0.01 [7]		
302	2.0481	0.1744	0.5857	2.2590	2.2544	
		0.1844	0.6190	2.2592	2.2543	
		0.1907	0.6400	2.2593	2.2543	
		0.2044	0.6855	2.2595	2.2541	
Expt.			<b>0.64</b> ± <b>0.01</b> [5]	<b>2.26</b> ± <b>0.02</b> [5]	$2.25{\pm}0.02*$	

Table 2. The ZFS parameters D and  $g_{\parallel}$ ,  $g_{\perp}$  for ZnSiF<sub>6</sub> · 6H<sub>2</sub>O:Ni<sup>2+</sup> at 4.2, 77 and 302 K.

* Obtained from NiSiF <sub>6</sub> · 6H <sub>2</sub> O in	[19].
---	-------

	4.2 K			77 K			298 K		
Level	Ca	ılc.	Obs.a	Ca	ılc.	Obs.a	Ca	ılc.	Obs.a
	A	В		A	В		A	$B^{b}$	
$^{3}T_{2}(^{3}F)$	9074	9161	9150	9043	9127	9120	8605	8796	8800
	9182	9172		9153	9143		8703	8845	
$^{3}T_{1}(^{3}F)$	15080	15173		15035	15124		14381	14646	14800
	15281	15193	15400	15240	15152	15290	14558	14735	
${}^{1}E({}^{1}D)$	15609	15612		15608	15610		15582	15593	
$^{1}T_{2}(^{1}D)$	24178	24333		24143	24293		23679	23916	
	24387	24353		24356	24322		23865	24008	
${}^{1}A_{1}({}^{1}G)$	25127	25124	24450	25117	25113		24942	24998	
$^{3}T_{1}(^{3}P)$	26277	26627	26100	26226	26564	26000	25594	25988	25650
	26804	26680		26763	26639		26067	26226	

Table 3. The observed and calculated optical spectral data of  $\mathrm{Ni}^{2+}$  in the two crystals  $\mathrm{NiTiF_6} \cdot \mathrm{6H_2O}$  (A) and  $\mathrm{ZnSiF_6} \cdot \mathrm{6H_2O}$  (B) at 4.2, 77 and 298 (302) K; N=0.969; all units in cm<sup>-1</sup>.

 $R = 2.0338 \text{ Å}, \Delta\theta = 0.4238^{\circ}; \text{ and } R = 2.0543 \text{ Å}, \Delta\theta =$  $0.3854^{\circ}$  for Ni<sup>2+</sup> in NiTiF<sub>6</sub> · 6H<sub>2</sub>O and R = 2.0327 Å,  $\Delta\theta = 0.0414^{\circ}$ ; R = 2.0341 Å,  $\Delta\theta = 0.0593^{\circ}$ ; and R =2.0481 Å,  $\Delta\theta = 0.1907^{\circ}$  for Ni<sup>2+</sup> in ZnSiF<sub>6</sub> · 6H<sub>2</sub>O are determined at 4.2, 77 and 298 (302) K, respectively. Simultaneously, we can determine the local structure parameters  $R = 2.0325 \text{ Å}, \ \theta = 55.1500^{\circ};$  $R = 2.0338 \text{ Å}, \ \theta = 55.1594^{\circ}; \text{ and } R = 2.0543 \text{ Å}, \ \theta =$ 55.1210° for Ni<sup>2+</sup> in NiTiF<sub>6</sub> · 6H<sub>2</sub>O and R = 2.0327 Å,  $\theta = 54.7770^{\circ}$ ; R = 2.0341 Å,  $\theta = 54.7949^{\circ}$ ; and R =2.0481 Å,  $\theta = 54.9263^{\circ}$  for Ni<sup>2+</sup> in ZnSiF<sub>6</sub> · 6H<sub>2</sub>O at temperatures 4.2, 77 and 298 (302) K, respectively. It must be pointed out that, due to the influence of  $\mathrm{H}^+$  ions, the calculated local structure parameters R and  $\theta$  are at the equivalent position to the effective charge of the  $H_2O$  molecule rather than that of the  $O^{2-}$ ion. According to the former theoretical and experimental researches,  $\theta < \theta_{\rm oh}$  represents the local lattice structure of the 3d<sup>n</sup> cations in trigonal crystals exhibiting an elongation distortion, contrarily,  $\theta > \theta_{\rm oh}$  represents the local lattice structure of the 3d<sup>n</sup> cations in

trigonal crystals exhibiting a compression distortion. The results of Tables 1 and 2 show that the local lattice structure around the Ni<sup>2+</sup> centre in NiTiF<sub>6</sub> · 6H<sub>2</sub>O and ZnSiF<sub>6</sub> · 6H<sub>2</sub>O has a compression distortion. Furthermore, the local lattice distortion degree for Ni<sup>2+</sup> in ZnSiF<sub>6</sub> · 6H<sub>2</sub>O increases with rising temperature, whereas for Ni<sup>2+</sup> in NiTiF<sub>6</sub> · 6H<sub>2</sub>O the angle distortion degree at 77 K is larger than at 4.2 K and 298 K. This exceptional variation for NiTiF<sub>6</sub> · 6H<sub>2</sub>O may be ascribed to the fact that the ZFS parameter D is very sensitive to the local structure parameter  $\theta$ , and the values of |D| observed in the experiment increase with temperature up to about 77 K, then decrease between about 77 K and 298 K.

### 4. Study of the Relationship between the Orbital Reduction Factor *k* and the *g* Factors

From (3), the orbital reduction factor k is taken into account for the Zeeman interaction in the perturbation Hamiltonian. By using the corresponding pa-

<sup>&</sup>lt;sup>a</sup> Spectral data obtained from [15].

<sup>&</sup>lt;sup>b</sup> At 302 K.

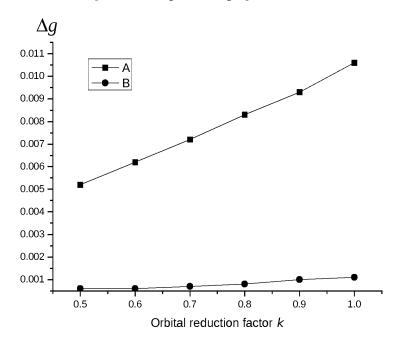


Fig. 2. The relationship between  $\Delta g$  and the orbital reduction factor k for NiTiF<sub>6</sub>·6H<sub>2</sub>O:Ni<sup>2+</sup> (A) and ZnSiF<sub>6</sub>·6H<sub>2</sub>O:Ni<sup>2+</sup> (B) at 4.2 K.

Table 4. The relationship between the EPR g factors and the orbital reduction factor k for NiTiF<sub>6</sub> · 6H<sub>2</sub>O:Ni<sup>2+</sup> at 4.2, 77 and 298 K.

-	4.2 K		77	' K	298 K	
k	$g_{\parallel}$	$g_{\perp}$	$g_{\parallel}$	$g_{\perp}$	$g_{\parallel}$	$g_{\perp}$
0.5	2.1262	2.1210	2.1266	2.1210	2.1321	2.1267
0.6	2.1528	2.1466	2.1534	2.1467	2.1601	2.1536
0.7	2.1794	2.1722	2.1801	2.1724	2.1881	2.1806
0.8	2.2061	2.1978	2.2068	2.1980	2.2160	2.2075
0.9	2.2327	2.2234	2.2335	2.2237	2.2440	2.2345
1.0	2.2593	2.2487	2.2603	2.2494	2.2720	2.2614

Table 5. The relationship between the EPR g factors and the orbital reduction factor k for ZnSiF<sub>6</sub> · 6H<sub>2</sub>O:Ni<sup>2+</sup> at 4.2, 77 and 302 K.

	4.2 K		77	' K	302 K	
k	$g_{\parallel}$	$g_{\perp}$	$g_{\parallel}$	$g_{\perp}$	$g_{\parallel}$	$g_{\perp}$
0.5	2.1230	2.1224	2.1235	2.1230	2.1285	2.1259
0.6	2.1490	2.1484	2.1497	2.1490	2.1558	2.1527
0.7	2.1750	2.1743	2.1758	2.1750	2.1830	2.1794
0.8	2.2010	2.2002	2.2019	2.2009	2.2103	2.2062
0.9	2.2271	2.2261	2.2281	2.2269	2.2375	2.2329
1.0	2.2531	2.2520	2.2542	2.2529	2.2648	2.2597

rameters at different temperatures above and adjusting the orbital reduction factor, the EPR parameters are determined depending on different k for Ni<sup>2+</sup> in NiTiF<sub>6</sub>·6H<sub>2</sub>O and ZnSiF<sub>6</sub>·6H<sub>2</sub>O, which are listed in Tables 4 and 5. Meanwhile, the relationship between  $\Delta g$  (=  $g_{\parallel} - g_{\perp}$ ) and the orbital reduction factor k for Ni<sup>2+</sup> in NiTiF<sub>6</sub>·6H<sub>2</sub>O and ZnSiF<sub>6</sub>·6H<sub>2</sub>O at

different temperatures is plotted in Figures 2-4. From Tables 4 and 5 we can see that both  $g_{\parallel}$  and  $g_{\perp}$  for the two systems increase obviously when the orbital reduction factor k increases. From Figs. 2-4, it is shown that: (i)  $\Delta g$  is positive at all values of k with a magnitude that increases monotonically with the orbital reduction factor k. Generally, the sign of  $\Delta g$  is related to the ZFS parameter D; for a  $d^8$  ion under the weak field approximation, the sign of  $\Delta g$  is contrary to that of the ZFS parameter D. Hence, for  $Ni^{2+}$  in  $NiTiF_6 \cdot 6H_2O$  and  $ZnSiF_6 \cdot 6H_2O$  crystals, the value of  $g_{\parallel}$  is larger than that of  $g_{\perp}$ , which corresponds to the negative ZFS parameter D. (ii) At each temperature, there is an almost linear relation between k and  $\Delta g$  for  $Ni^{2+}$  in  $NiTiF_6 \cdot 6H_2O$  and  $ZnSiF_6 \cdot 6H_2O$ .

### 5. Conclusion

By diagonalizing the complete energy matrices, the local structure parameters and the zero-field splitting parameter D as well as the g factors of NiTiF<sub>6</sub> · 6H<sub>2</sub>O and ZnSiF<sub>6</sub> · 6H<sub>2</sub>O:Ni<sup>2+</sup>have been studied at different temperatures. It has been shown that the local lattice structure around the octahedral Ni<sup>2+</sup> centre in NiTiF<sub>6</sub> · 6H<sub>2</sub>O and ZnSiF<sub>6</sub> · 6H<sub>2</sub>O exhibits a compression distortion. The angle distortion degree of NiTiF<sub>6</sub> · 6H<sub>2</sub>O at 77 K is larger than that at 4.2 K and 298 K. For the ZnSiF<sub>6</sub> · 6H<sub>2</sub>O:Ni<sup>2+</sup> system, the compression distortion may be ascribed to the fact that

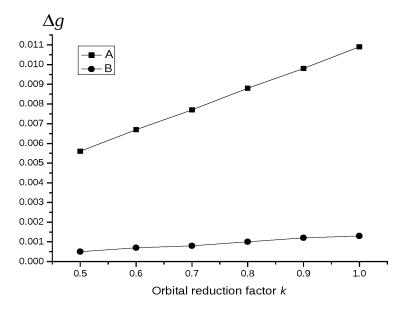


Fig. 3. The relationship between  $\Delta g$  and the orbital reduction factor k for NiTiF<sub>6</sub> · 6H<sub>2</sub>O:Ni<sup>2+</sup> (A) and ZnSiF<sub>6</sub> · 6H<sub>2</sub>O:Ni<sup>2+</sup> (B) at 77 K.

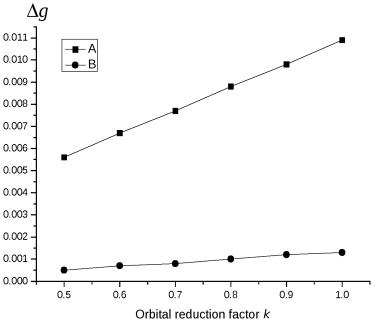


Fig. 4. The relationship between  $\Delta g$  and the orbital reduction factor k for NiTiF<sub>6</sub>·6H<sub>2</sub>O:Ni<sup>2+</sup> (A) and ZnSiF<sub>6</sub>·6H<sub>2</sub>O:Ni<sup>2+</sup> (B) at room temperature.

the radius of the Ni<sup>2+</sup> ion (0.72 Å) is smaller than that of the host ion Zn<sup>2+</sup> (0.75 Å) and the local lattice distortion degree of Ni<sup>2+</sup> in ZnSiF<sub>6</sub> · 6H<sub>2</sub>O increases with rising temperature. In addition, from our calculation, we also found that the EPR g factors  $g_{\parallel}$  and  $g_{\perp}$  depend sensitively on the orbital reduction factor k and there is a nearly linear relation between k and  $\Delta g$  for Ni<sup>2+</sup> in NiTiF<sub>6</sub> · 6H<sub>2</sub>O and ZnSiF<sub>6</sub> · 6H<sub>2</sub>O at each temperature.

### Acknowledgement

This work was supported by the National Natural Science Foundation of China (No. 10774103) and the Doctoral Education Fund of the Education Ministry of China (No. 20050610011).

- [1] L. Pauling, Z. Kristallogr. 72, 482 (1930).
- [2] R. S. Rubins, B. C. Griffin, and R. Burris, J. Chem. Phys. 64, 3349 (1976).
- [3] S. A. Al'tshuler and R. M. Valishev, Sov. Phys. JETP 21, 309 (1965).
- [4] R. S. Rubins, J. Chem. Phys. 84, 4142 (1986).
- [5] R. S. Rubins and S. K. Jani, J. Chem. Phys. 66, 3297 (1977).
- [6] R. H. Hoskins, R. C. Pastor, and K. R. Trigger, J. Chem. Phys. 30, 601 (1959).
- [7] I. J. Fritz and L. Yarmus, J. Chem. Phys. 51, 1428 (1969).
- [8] A. Abragam and B. Bleaney, Electron Paramagnetic Resonance of Transition Ions, Oxford University Press, Oxford 1970.
- [9] X. Y. Kuang, Phys. Rev. B 36, 712 (1987); B 36, 797 (1987); X. Y. Kuang, Ph. D. Thesis, Université Parissud, 1994, p. 14.

- [10] O. Kahn, Molecular Magnetism, VCH, New York 1993, p. 31.
- [11] D. J. Newman and W. Urban, Adv. Phys. 24, 793 (1975).
- [12] J. H. VanVleck, J. Chem. Phys. 1, 208 (1932).
- [13] M. G. Zhao, M. L. Du, and G. Y. Sen, J. Phys. C: Solid State Phys. 20, 5557 (1987).
- [14] S. Ray, A. Zalkin, and D. H. Templeton, Acta Cryst. B 29, 2741 (1973).
- [15] M. H. L. Pryce, G. Agnetta, T. Garofano, M. B. Paima-Vittorelli, and M. U. Palma, Phil. Mag. 10, 477 (1964).
- [16] D. Curie, C. Barthon, and B. Canny, J. Chem. Phys. 61, 3048 (1974).
- [17] J. S. Griffith, The Theory of Transition-Metal Ions, Cambridge University Press, London 1964.
- [18] W. C. Zheng, Radiat. Eff. Defect Solids 140, 329 (1997).
- [19] R.S. Rubins, J.D. Clark, and S.K. Jani, J. Chem. Phys. 67, 893 (1977).